Looking for Werner at RHIC

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easy

- News
- Publication talks
- Proposals

Brookhaven National Laboratory Nuclear and Particle Physics Program Advisory Committee January 2010

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particles that has not been adequately answered, despite decades of study.

Looking Forward (rapidity) at RHIC

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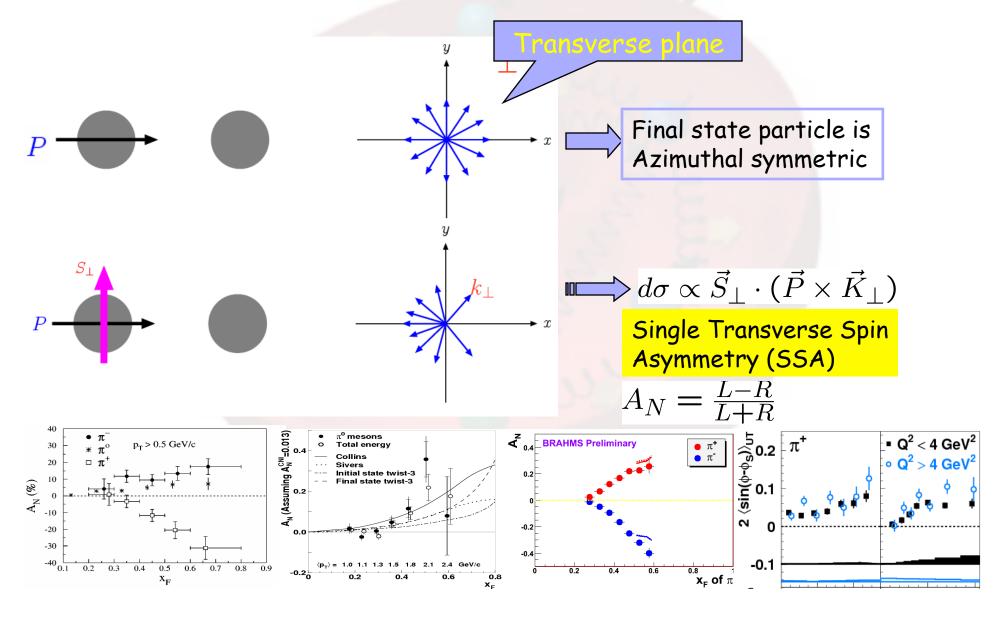
Great opportunities at forward direction of RHIC

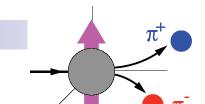
- Transverse spin phenomena
- Saturation physics in dA collisions
- Many others...





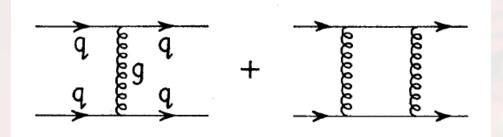
Single spin asymmetry at forward





Naïve parton model fails

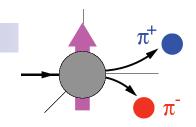
■ If the underlying scattering mechanism is hard, the naïve parton model generates a very small SSA: (G. Kane et al, 1978),



- \Box It is in general suppressed by $\alpha_s m_q/Q$
- We have to go beyond this naïve picture

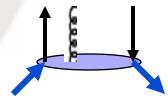






Two mechanisms in QCD

- Spin-dependent transverse momentum dependent (TMD) function s_{TA}
 - □ Sivers 90
- Sivers function $\sim S_T (PXk_T)$
- □ Colllins 93
- □ Brodsky-Hwang-Schmidt, 02
- Twist-3 quark-gluon correlations (coll.)
 - □ Efremov-Teryaev, 82, 84
 - □ Qiu-Sterman, 91,98







Single transverse-spin asymmetry in high transverse momentum pion production in pp collisions

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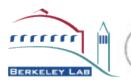
(Received 25 September 2006; revised manuscript received 10 November 2006; published 15 December 2006)

We study the single-spin (left-right) asymmetry in single-inclusive pion production in hadronic scattering. This asymmetry is power-suppressed in the transverse momentum of the produced pion and can be analyzed in terms of twist-three parton correlation functions in the proton. We present new calculations of the corresponding partonic hard-scattering functions that include the so-called "non-derivative" contributions not previously considered in the literature. We find a remarkably simple structure of the results. We also present a brief phenomenological study of the spin asymmetry, taking into account data from fixed-target scattering and also the latest information available from Relativistic Heavy Ion Collider (RHIC). We make additional predictions that may be tested experimentally at RHIC.

 $F = \gamma$, jet, pion, W, ...

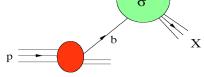
$$\begin{split} \Delta\sigma_{A+B\to hX}(\ell_{\perp},\vec{s}_{T}) &= \sum_{abc} \phi_{a/A}^{(3)}(x_{1},x_{2},\vec{s}_{T}) \otimes \phi_{b/B}(x') \otimes H_{ab\to c}(\ell_{\perp},\vec{s}_{T}) \otimes D_{c\to h}(z) \\ \textbf{Qiu-Sterman} &+ \sum_{abc} \delta q_{a/A}(x,\vec{s}_{T}) \otimes \phi_{b/B}^{(3)}(x'_{1},x'_{2}) \otimes H'_{ab\to c}(\ell_{\perp},\vec{s}_{T}) \otimes D_{c\to h}(z) \\ &+ \sum_{abc} \delta q_{a/A}(x,\vec{s}_{T}) \otimes \phi_{b/B}(x') \otimes H''_{ab\to c}(\ell_{\perp},\vec{s}_{T}) \otimes D_{c\to h}^{(3)}(z_{1},z_{2}) \end{split}$$

Qiu-Sterman;KQVY 06,10 Kanazawa-Koike 00 Kang-Yuan-Zhou 10

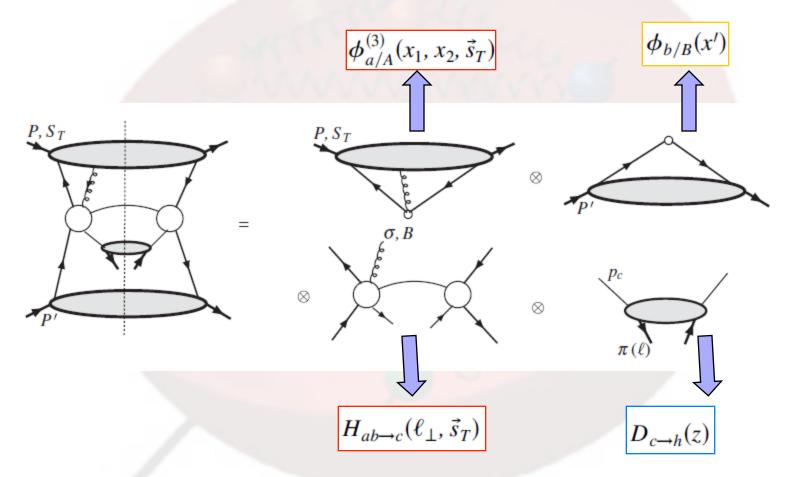








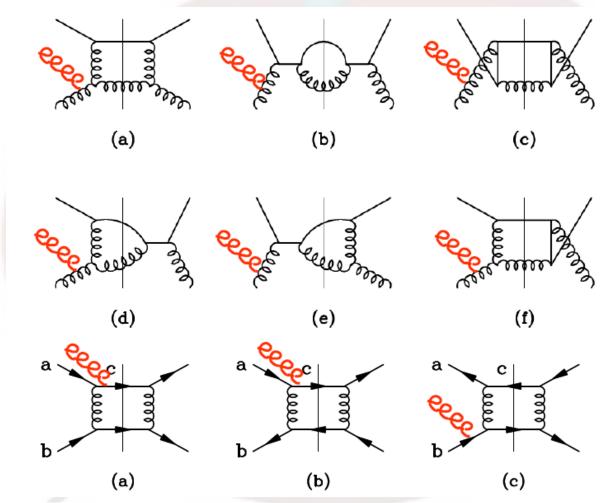
Collinear factorization







Twist-three diagrams







Remarkable simple results

 single inclusive hadron production in pp collision, including the derivative and non-derivative terms

$$E_{\ell} \frac{d^{3} \Delta \sigma(\vec{s}_{T})}{d^{3} \ell} = \frac{\alpha_{s}^{2}}{S} \sum_{a,b,c} \int_{z_{\min}}^{1} \frac{dz}{z^{2}} D_{c \to h}(z) \int_{x'_{\min}}^{1} \frac{dx'}{x'} \frac{1}{x'S + T/z} \phi_{b/B}(x')$$

$$\times \sqrt{4\pi \alpha_{s}} \left(\frac{\epsilon^{\ell s_{T} n\bar{n}}}{z\hat{u}} \right) \frac{1}{x} \left[T_{a,F}(x,x) - x \left(\frac{d}{dx} T_{a,F}(x,x) \right) \right] H_{ab \to c}(\hat{s}, \hat{t}, \hat{u})$$

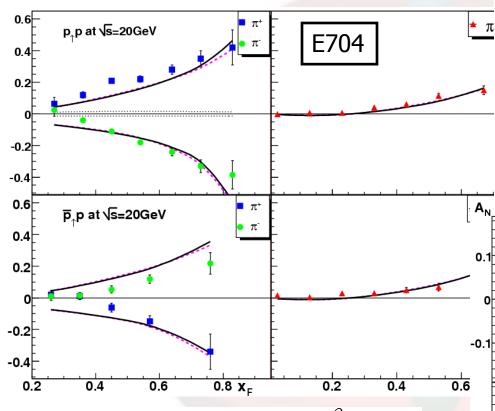
Qiu, Sterman, 91, 98 Kouvaris, Qiu, Vogelsang, Yuan, 06





Twist-3 Fit to data

$$p_{\uparrow}p \to \pi + X$$

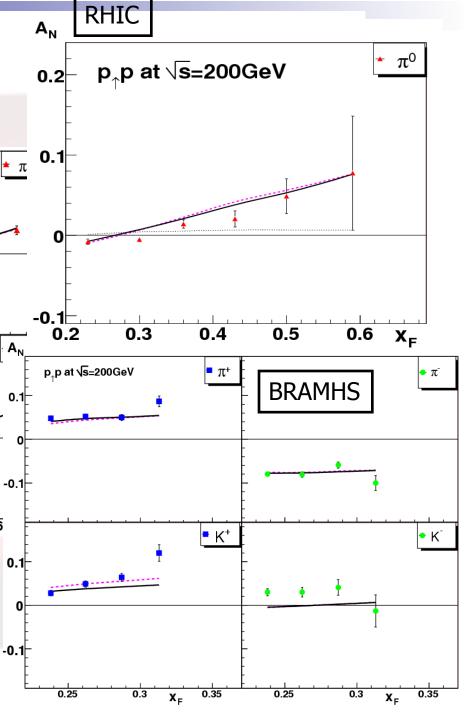


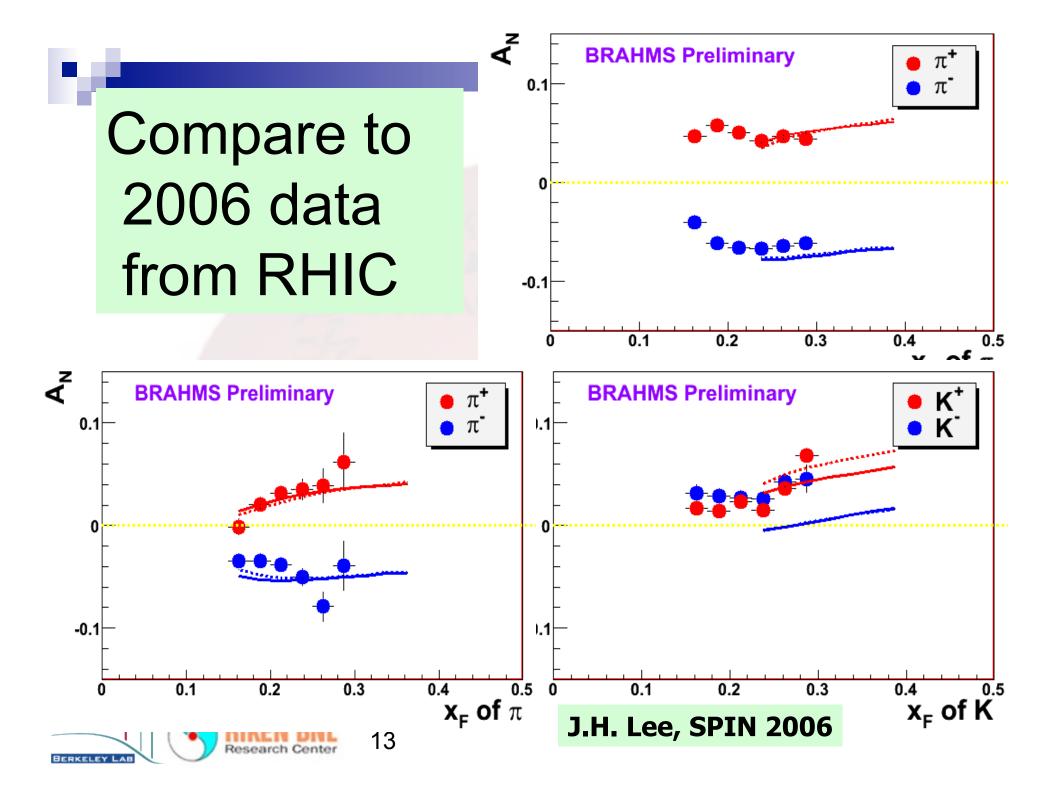
$$T_F^a(x) = N_a x^{\alpha_a} (1 - x)^{\beta_a} f_a(x)$$

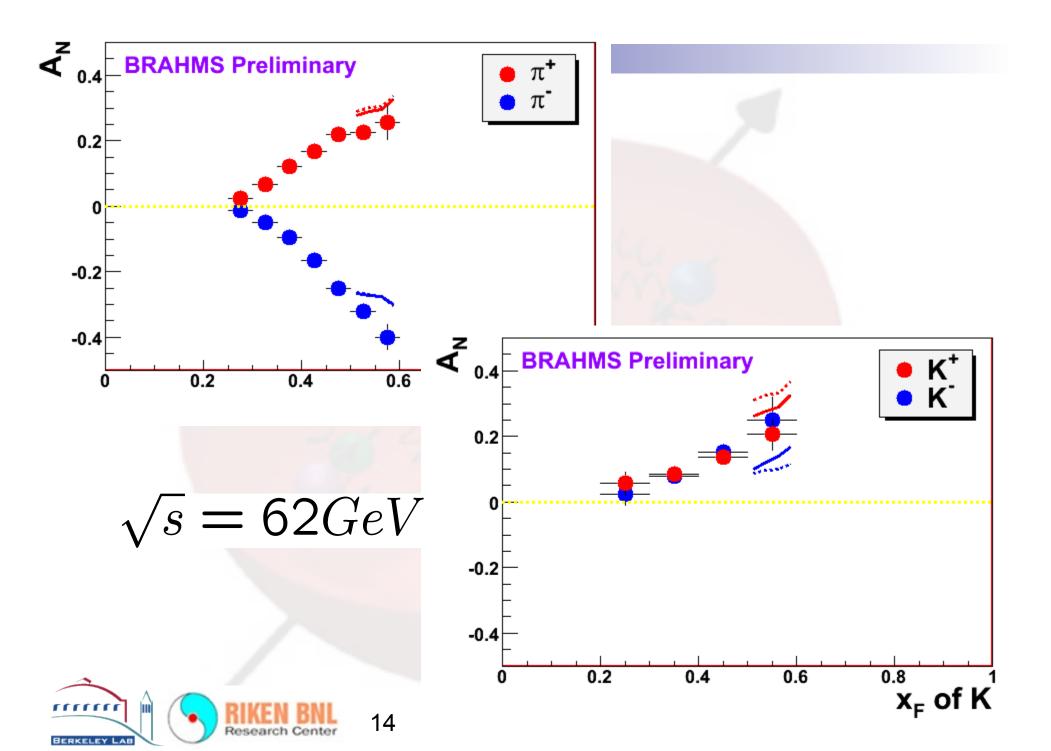
Kouvaris, Qiu, Vogelsang, Yuan, 06













We are also updating

- Gluon contribution
- Collins (twist-three) fragmentation contribution
- Stay tune,...

Kang, Qiu, Vogelsang, Yuan, to appear





NLO corrections to SSA

Vogelsang-Yuan, arXiv:0904.0410

■ SSA in Drell-Yan as an example,

$$p_{\uparrow}(P_A, S_{\perp}) p(P_B) \rightarrow \gamma^*(Q^2, q_{\perp}) + X \rightarrow \ell^+\ell^- + X$$

$$|\epsilon^{\alpha\beta}S_{\perp}^{\alpha}q_{\perp}^{\beta}| = |S_{\perp}||q_{\perp}|\sin\phi$$

■ Collinear factorization

3/22/10

$$\frac{d\langle q_{\perp} \Delta \sigma(S_{\perp}) \rangle}{dQ^2} = \sigma_0 \int \frac{dx_1}{x_1} \frac{dx_2}{x_2} \frac{dx'}{x'} T_{F,q}(x_1, x_2) \bar{q}(x') \mathcal{H}(x_1, x_2; x')$$

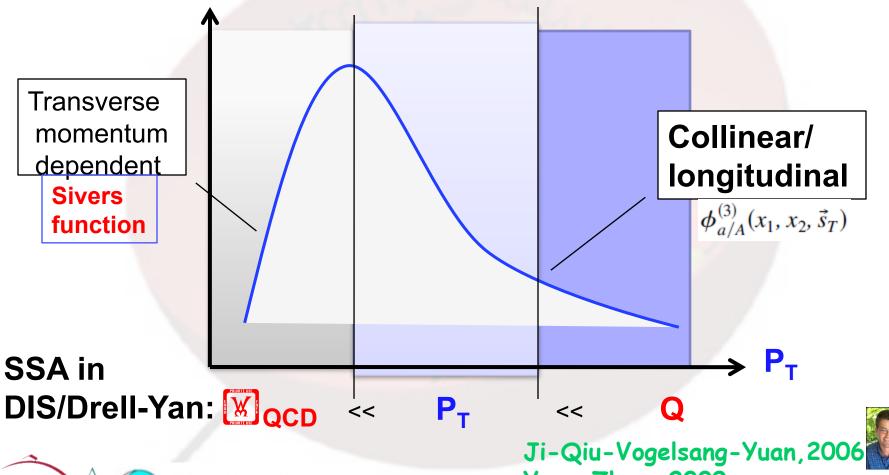
Collinear functions, evolution: Kang-Qiu, 08;

Zhou-Yuan-Liang, 08 Braun et al., 0909.3410





A unified picture between TMD and twist-three (leading pt/Q)



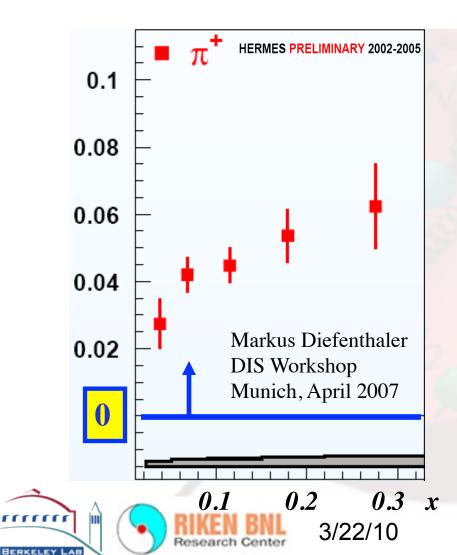


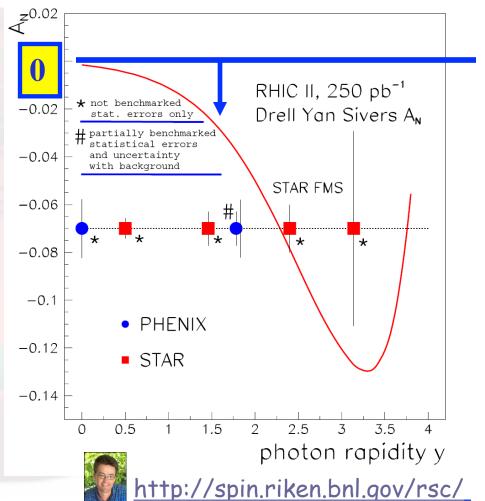


Experiment SIDIS vs Drell Yan

HERMES Sivers Results

RHIC II Drell Yan Projections

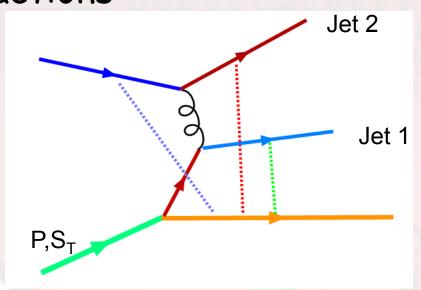




<u>les Bland, preliminary</u>

Non-universality: Dijet-correlation at RHIC

Initial state and/or final state interactions





Standard Factorization breaks, no universality!



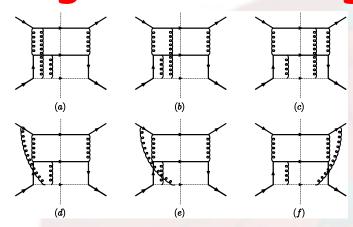




The simple picture does not hold for Vogelang-Yuan, 0708.4398;

two-gluon exchanges



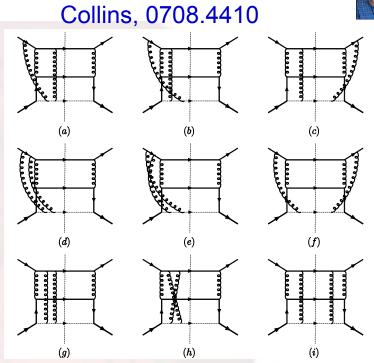


Becchetta-Bomhof-Mulders-Pijlman, 04-06

$$\mathcal{L}'_{v_a}(g_1, g_2; \xi) \equiv \mathcal{P} \exp\left(-ig_1 \int_0^\infty d\lambda v_a \cdot A(\xi + \lambda v_a)\right)$$

$$\times \mathcal{P} \exp\left(-ig_2 \int_0^\infty d\lambda v_a \cdot A(\xi + \lambda v_a)\right)$$

$$\times \mathcal{P} \exp\left(ig_2 \int_0^{-\infty} d\lambda v_a \cdot A(\xi + \lambda v_a)\right)$$



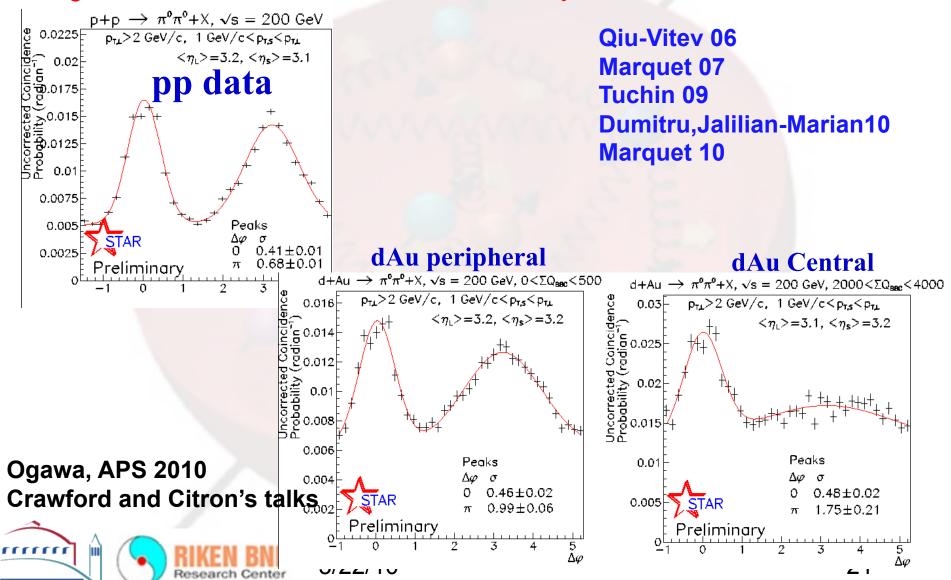
$$\mathcal{P}\exp\left(-ig_1\int_0^\infty d\lambda v_a\cdot A(\xi+\lambda v_a)\right)$$

Integrated over transverse momentum





Dijet-correlation in pA collisions



NA.

Universality of TMD at small-x?

- This has not been shown/studied in small-x physics
- Different assumption and summation are made
- We want to make a model calculations suitable for both TMD and small-x approximation
 - Summation to all order is crucial

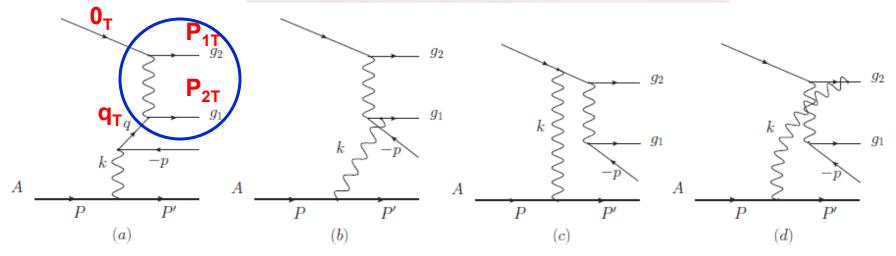






 $qq' \rightarrow qq'$ channel

Xiao, Yuan, arXiv:1003.0482



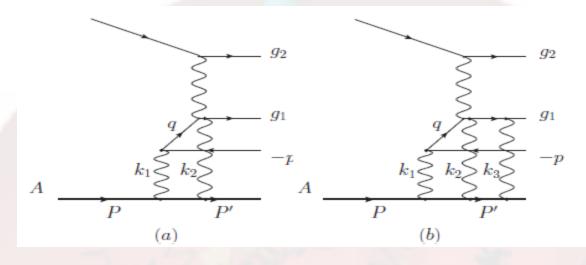
 $= q_t \text{ dependence, } q_T = P_{1T} + P_{1T} \leftrightarrow P_{1T} \sim P_{1T}$

$$\tilde{q}\left(x,q_{\perp}\right) = \frac{x}{32\pi^{2}} \int \frac{dp^{-}}{p^{-}} \frac{d^{2}k_{\perp}}{\left(2\pi\right)^{4}} (4P^{+}p^{-})^{2} \left|A^{(tot)}\left(k,p\right)\right|^{2} \\ A^{(1)}\left(k,p\right) = gg_{1} \frac{1}{k_{\perp}^{2} + \lambda^{2}} \left[\frac{1}{D_{1}} - \frac{1}{D_{2}}\right]$$





Multi-gluon exchange



$$\begin{split} A^{(2)}\left(k,p\right) &= \frac{i}{2}g^2 \int d[1]d[2] \left\{ g_1^2 \left[\frac{1}{D_1} + \frac{1}{D_2} - \frac{1}{D_{21}} - \frac{1}{D_{22}} \right] + g_1g_2 \left[\frac{2}{D_2} - \frac{2}{D_{21}} \right] \right\} \\ A^{(3)}\left(k,p\right) &= \frac{1}{3!}g^3 \int d[1]d[2]d[3] \left\{ g_1^3 \left[\frac{1}{D_2} - \frac{1}{D_1} + \frac{3}{D_{13}} - \frac{3}{D_{21}} \right] + g_1g_2^2 \left[\frac{3}{D_2} - \frac{3}{D_{21}} \right] \right\} \\ &+ g_1^2g_2 \left[\frac{3}{D_2} + \frac{3}{D_{13}} - \frac{3}{D_{21}} - \frac{3}{D_{22}} \right] + g_1g_2^2 \left[\frac{3}{D_2} - \frac{3}{D_{21}} \right] \right\} \end{split}$$





All-orders

$$\begin{split} \tilde{q} \left(x, q_{\perp} \right) &= \frac{x P^{+2}}{8 \pi^4} \int dp^- p^- \int d^2 R_{\perp} d^2 R_{\perp}' d^2 r_{\perp} e^{i q_{\perp} \cdot \left(R_{\perp} - R_{\perp}' \right)} e^{-i g g_2 \left(G(R_{\perp}) - G(R_{\perp}') \right)} \\ V \left(r_{\perp} \right) V \left(r_{\perp}' \right) &\left\{ 1 - e^{i g g_1 \left[G\left(R_{\perp} + r_{\perp} \right) - G(R_{\perp}) \right]} \right\} \left\{ 1 - e^{-i g g_1 \left[G\left(R_{\perp}' + r_{\perp}' \right) - G\left(R_{\perp}' \right) \right]} \right\} \; , \end{split}$$

■ DIS process $(g_2$ -term disappear)

$$\begin{split} \tilde{q} \left(x, q_{\perp} \right) &= \frac{x P^{+2}}{8 \pi^4} \int dp^- p^- \int d^2 R_{\perp} d^2 R_{\perp}' d^2 r_{\perp} e^{i q_{\perp} \cdot \left(R_{\perp} - R_{\perp}' \right)} V \left(r_{\perp} \right) V \left(r_{\perp}' \right) \\ &\times \left\{ 1 - e^{i g g_1 \left[G \left(R_{\perp} + r_{\perp} \right) - G \left(R_{\perp} \right) \right]} \right\} \left\{ 1 - e^{-i g g_1 \left[G \left(R_{\perp}' + r_{\perp}' \right) - G \left(R_{\perp}' \right) \right]} \right\} \; , \end{split}$$

■ They are not the same, Non-universality

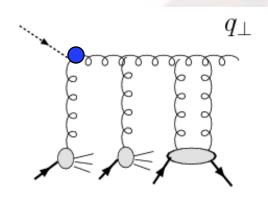


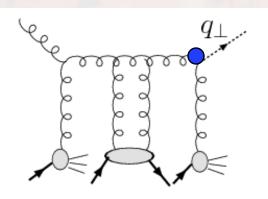


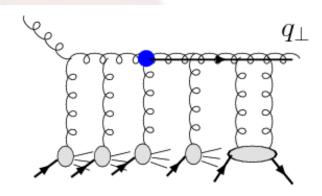


Realistic QCD dipole model

Following Kovchegov-Mueller 1998









Drell-Yan-type

$$\tilde{N}(\underline{x}) = \int d^2b \frac{N_c^2 - 1}{\pi^2 \alpha N_c \underline{x}^2} (1 - exp[-\frac{\sqrt{R^2 - b^2}}{2\lambda} \underline{x}^2 \tilde{v}(\underline{x})])$$

Time-reversal invariance

pA→Dijet-like



Certainly will be different

Marquet, Venugopalan, Xiao, Yuan, work in progress







Comments

- Light-cone gauge does not help
- Non-universal for the kt-dependent parton distribution at small-x, will affect the phenomenological interpretations
- Remain to be seen that how/does the classical field calculation contain these effects





100

Journey forward at RHIC

- SSA in forward direction impose theoretical challenge
 - □Pt-dependence, Kang et al, work in progress
 - □Eta/piO SSA
- Dijet-correlation not only probe small-x saturation, but also the QCD dynamics (initial/final state interaction effects)



